#### Notes:

# Exam:

- Average was 46.
- You can make up some points:
  - 1. Choose 2 problems. I will post the blank exam.
  - 2. Work out the solutions from scratch on your own.
  - 3. Turn in the problems by Thursday in class or Friday afternoon (Tentatively 3 PM, I will arrange to be at my office). Not in my mailbox. Full details to come.
  - 4. You will get 60% of the made-up points back.
  - 5. You can look for a few minutes at the exam now but not take it with you.

**Homework:** Note problem set due Thursday.

## **Canonical ensemble Recall:**

Ζ

Partition function

$$= \sum_{states i} Exp[-E_i/kT] \qquad F \equiv -k_B T \ln Z$$

$$dF = -S dT - P dV + \mu N$$

$$S = -\left(\frac{\partial F}{\partial T}\right)_{V,N} \qquad P = -\left(\frac{\partial F}{\partial V}\right)_{T,N} \qquad \mu = \left(\frac{\partial F}{\partial N}\right)_{T,V}$$

$$S = k_B \ln Z + k_B T \frac{\partial}{\partial T} \ln Z$$

$$\langle E \rangle = k_B T^2 \frac{\partial}{\partial T} \ln Z$$

#### More on entropy:

$$F \equiv -k_B T ln Z$$

- Free energy can treat as defined quantity.
- *F* is a conserved quantity, doesn't fluctuate.
- Equivalent to thermodynamic F, we will see.

Then: 
$$S = -\left(\frac{\partial F}{\partial T}\right)_{V,N}$$
  $(dF = -SdT - PdV + \mu N)$ 

From here can show:  

$$S = -k_B \sum_{j} P_j \ln(P_j)$$

<u>Gibbs</u> version of entropy, vs. <u>Boltzmann</u> version appropriate for Microcanonical:  $S = k_B \ln(\Omega)$ 

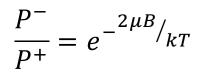
Note this form of entropy is in chapter 17. Figures prominently in <u>information theory</u>.

# spin-1/2 non-interacting paramagnet

$$E = \pm \mu B$$
 per atom ->  $E = \mu B(N^- - N^+) = \mu B(2N^- - N)$ 

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N distinguishable atoms



Also from last time: Z as *product* of individual components. Advantageous solution method.

 $Z_{i} = e^{+\frac{\mu B}{kT} + e^{-\frac{\mu B}{kT}}}$ (paramagnet specific case)  $= 2 \cosh^{\frac{\mu B}{kT}} = 2 \cosh^{\frac{\mu B}{kT}}$  $Z = \prod_{i} Z_{i}; \text{ includes all cross terms.}$ 

Works in general if:

- Energies add (classical *separation of variables*; quantum states as simple product wavefunction);
- Occupation probabilities independent of each other.

#### **Equipartition theorem**

- <u>Classical</u> systems (continuum not discrete energies)
- Works in cases having separable variables.
- Requires <u>energy quadratic in position and/or momentum:</u>

 $E = cq^2$  Text notation (means variables q and p).

Result: for <u>any</u> energy of this general form, easy to show,

$$U = \frac{f}{2}kT$$

from canonical partition function, (1/2)kT obtained for each "degree of freedom" *f*.

$$Z = \frac{1}{h^{3N}} \iiint d^{3N} r d^{3N} p e^{-\beta E} - Z = \prod_i Z_i; \text{ includes all cross terms.}$$

Classical partition function

systems of distinguishable particles, non-interacting.

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Classical partition function

systems of distinguishable particles, non-interacting.

$$Z = \frac{1}{N!} \prod_{i} Z_i$$

systems of **indistinguishable** particles, still <u>non-interacting</u> case.

#### Harmonic oscillator systems

<u>Classical</u> systems: 

$$U_i = \frac{\overline{p_u}^2}{2m} + \frac{\kappa \overline{u}^2}{2} \qquad \text{or} \qquad U_i = \frac{p^2}{2m} + \frac{\kappa u^2}{2}$$

c.m. coordinates & reduced mass

Classical partition function

$$Z = \frac{1}{N! h^{3N}} \iiint d^{3N} r d^{3N} p e^{-\beta E_{trans}} \left(\frac{1}{h} \iint dp du e^{-\beta \left(\frac{p^2}{2m} + \frac{\kappa u^2}{2}\right)}\right)^N$$
$$\langle E \rangle = \frac{3}{n} NkT$$
Internal degrees of freedom

$$E\rangle = \frac{3}{2}NkT$$

translational

Internal degrees of freedom (rotational more important for e.g.  $N_2$  at room temperature as we have seen)

## Harmonic oscillator systems

• <u>Classical</u> systems:

$$U_i = \frac{\overline{p_u}^2}{2m} + \frac{\kappa \overline{u}^2}{2}$$
 or  $U_i = \frac{p^2}{2m} + \frac{\kappa u^2}{2}$ 

c.m. coordinates & reduced mass

Classical partition function – harmonic-potential terms only:

$$Z = \left(\frac{1}{h} \iint dp du e^{-\beta \left(\frac{p^2}{2m} + \frac{\kappa u^2}{2}\right)}\right)^N$$

N independent oscillators (Einstein solid approximation, ideal gas molecules, etc.)

# Equipartition theorem:

- <u>Classical</u> systems (continuum not discrete energies)
- Works in cases having separable variables.
- Requires <u>energy quadratic in position and/or momentum</u> (or other coordinate)
- Vibrations, rotations, translational states. Rotational case in text, won't show here.