

# Complete characterization of fourth-order symplectic integrators with extended-linear coefficients

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The structure of symplectic integrators up to fourth order can be completely and analytically understood when the factorization (split) coefficients are related linearly but with a uniform nonlinear proportional factor. The analytic form of these *extended-linear* symplectic integrators greatly simplified proofs of their general properties and allowed easy construction of both forward and nonforward fourth-order algorithms with an arbitrary number of operators. Most fourth-order forward integrators can now be derived analytically from this extended-linear formulation without the use of symbolic algebra.

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## I. INTRODUCTION

Evolution equations of the form

$$w(t + \varepsilon) = e^{\varepsilon(T+V)}w(t), \quad (1.1)$$

where  $T$  and  $V$  are noncommuting operators, are fundamental to all fields of physics ranging from classical mechanics [1–5], electrodynamics [6,7], and statistical mechanics [8–11] to quantum mechanics [12–14]. All can be solved by approximating  $e^{\varepsilon(T+V)}$  to the  $(n+1)$ th order in the product form

$$e^{\varepsilon(T+V)} = \prod_{i=1}^N e^{t_i \varepsilon T} e^{v_i \varepsilon V} + O(\varepsilon^{n+1}) \quad (1.2)$$

via a well chosen set of factorization (or split) coefficients  $\{t_i\}$  and  $\{v_i\}$ . The resulting algorithm is then  $n$ th order because the algorithm's Hamiltonian is  $T+V+O(\varepsilon^n)$ . By understanding this single approximation, computational problems in diverse fields of physics can all be solved by applying the same algorithm.

Classically, (1.2) results in a class of composed or factorized symplectic integrators. While the conditions on  $\{t_i\}$  and  $\{v_i\}$  for producing an  $n$ th order algorithm can be stated, these *order conditions* are highly nonlinear and analytically opaque. In many cases [14–17], elaborate symbolic mathematical programs are needed to produce even fairly low order algorithms if  $N$  is large. In this work, we show that the structure of most fourth-order algorithms, including nearly all known forward ( $\{t_i, v_i\} > 0$ ) integrators, can be understood and derived on the basis that  $\{v_i\}$  and  $\{t_i\}$  are linearly related but with a uniform nonlinear proportional factor. This class of *extended-linear integrators* is sufficiently complex to be representative of symplectic algorithms in general, but its transparent structure makes it invaluable for constructing integrators up to the fourth-order. In this work we prove three important theorems, on the basis of which many families of fourth-order algorithms can be derived with analytically known coefficients, including all known forward integrators up to  $N=4$ .

## II. THE ERROR COEFFICIENTS

The product form (1.2) has the general expansion

$$\prod_{i=1}^N e^{t_i \varepsilon T} e^{v_i \varepsilon V} = \exp\{\varepsilon e_T T + \varepsilon e_V V + \varepsilon^2 e_{TV}[T, V] + \varepsilon^3 e_{TTV}[T, [T, V]] + \varepsilon^3 e_{VTV}[V, [T, V]] + \dots\}. \quad (2.1)$$

We have previously [18] described in detail how the error coefficients  $e_T$ ,  $e_V$ ,  $e_{TV}$ ,  $e_{TTV}$ , and  $e_{VTV}$  can be determined from  $\{t_i\}$  and  $\{v_i\}$ :

$$e_T = \sum_{i=1}^N t_i, \quad e_V = \sum_{i=1}^N v_i, \quad (2.2)$$

$$\frac{1}{2} + e_{TV} = \sum_{i=1}^N \nabla s_i u_i, \quad (2.3)$$

$$\frac{1}{3!} + \frac{1}{2} e_{TV} + e_{TTV} = \frac{1}{2} \sum_{i=1}^N \nabla s_i^2 u_i, \quad (2.4)$$

$$\frac{1}{3!} + \frac{1}{2} e_{TV} - e_{VTV} = \frac{1}{2} \sum_{i=1}^N \nabla s_i u_i^2, \quad (2.5)$$

where we have defined useful variables

$$s_i = \sum_{j=1}^i t_j, \quad u_i = \sum_{j=i}^N v_j, \quad (2.6)$$

with  $s_0=0$ ,  $u_{N+1}=0$ , and the backward finite differences

$$\nabla s_i^n = s_i^n - s_{i-1}^n, \quad (2.7)$$

with property

$$\sum_{i=1}^N \nabla s_i^n = s_N^n (= e_T^n = 1). \quad (2.8)$$

We will always assume that the primary constraint  $e_T=1$  and  $e_V=1$  are satisfied so that (2.8) sums to unity. Satisfying these two primary constraints is sufficient to produce a first-order algorithm. For a second-order algorithm, one must additionally force  $e_{TV}=0$ . For a third-order algorithm, one further requires that  $e_{TTV}=0$  and  $e_{VTV}=0$ . For a fourth-order

algorithm, it is sufficient to satisfy the third-order constraints with coefficients  $t_i$  that are left-right symmetric. (The symmetry for  $v_i$  will follow and need not be imposed *a priori*.) Once the primary conditions  $e_T=1$  and  $e_V=1$  are imposed, the constraints equations (2.3)–(2.5) are highly nonlinear and difficult to decipher analytically. In this work, we will show that (2.3) can be satisfied for all  $N$  by having  $\{v_i\}$  linearly related to  $\{t_i\}$  (or vice versa). The coefficients  $e_{TTV}$  and  $e_{VTV}$  can then be evaluated simply in terms of  $\{t_i\}$  (or  $\{v_i\}$ ) alone. This then completely determines the structure of third- and fourth-order algorithms.

### III. THE EXTENDED-LINEAR FORMULATION

The constraint  $e_{TV}=0$  is satisfied if

$$\sum_{i=1}^N \nabla s_i u_i = \frac{1}{2}. \quad (3.1)$$

If we view  $\{t_i\}$  as given, this is a linear equation for  $\{u_i\}$ . Knowing (2.8), a general solution for  $u_i$  in terms of  $s_i$  and  $s_{i-1}$  is

$$u_i = \sum_{n=1}^M C_n \frac{\nabla s_i^n}{\nabla s_i}, \quad \text{with } \sum_{n=1}^M C_n = \frac{1}{2}. \quad (3.2)$$

The coefficients  $C_n$  represent the intrinsic freedom in  $\{v_i\}$  to satisfy any constraint as expressed through its relationship to  $\{t_i\}$ . The expansion (3.2) is in increasing powers of  $s_i$  and  $s_{i-1}$ . If we truncated the expansion at  $M=2$ , then for  $i \neq 1$ ,  $u_i$  is linearly related to  $\{s_i\}$ , i.e.,

$$u_i = C_1 + C_2 \frac{\nabla s_i^2}{\nabla s_i} = C_1 + C_2(s_i + s_{i-1}). \quad (3.3)$$

For  $i=1$ , since we must satisfy the primary constraint  $e_V=1$ , we must have

$$u_1 = 1. \quad (3.4)$$

In this case, the constraint (3.1) takes the form

$$\sum_{i=1}^N \nabla s_i u_i = t_1 + C_1(1 - t_1) + C_2(1 - t_1^2) = \frac{1}{2}. \quad (3.5)$$

The complication introduced by  $u_1=1$ , in this, and in other similar sums, can be avoided without any loss of generality by decreeing

$$t_1 = 0, \quad (3.6)$$

so that (3.5) remains

$$C_1 + C_2 = \frac{1}{2}. \quad (3.7)$$

For  $i \neq 1 \neq N$ , (3.3) implies that

$$v_i = -C_2(t_i + t_{i+1}). \quad (3.8)$$

Since  $v_1 = u_1 - u_2 = 1 - C_1 - C_2 t_2$ , by virtue of (3.7),

$$v_1 = \frac{1}{2} + C_2(1 - t_2). \quad (3.9)$$

Similarly, since  $v_N = u_N = C_1 + C_2(2 - t_N)$ , we also have

$$v_N = \frac{1}{2} + C_2(1 - t_N). \quad (3.10)$$

Given  $\{t_i\}$  such that  $t_1=0$ , the set of  $\{v_i\}$  defined by (3.8)–(3.10) automatically satisfies  $e_V=1$  and  $e_{TV}=0$ . If  $C_2$  were a real constant, then  $\{v_i\}$  is linearly related to  $\{t_i\}$ . However, in most cases  $C_2$  will be a function of  $\{t_i\}$  and the actual dependence is nonlinear. But the nonlinearity is restricted to  $C_2$ , which is the same for all  $v_i$ . We will call this special form of dependence of  $v_i$  on  $\{t_i\}$ , *extended-linear*. For a given set of  $t_i$ , (3.8)–(3.10) defines our class of extended-linear integrators with one remaining parameter  $C_2$ .

For extended-linear integrators as described above, one can easily check that the sums in (2.4) and (2.5) can be evaluated as

$$\sum_{i=1}^N \nabla s_i^2 u_i = C_1 + C_2 + g C_2 = \frac{1}{2} + g C_2, \quad (3.11)$$

$$\sum_{i=1}^N \nabla s_i u_i^2 = (C_1 + C_2)^2 + g C_2^2 = \frac{1}{4} + g C_2^2. \quad (3.12)$$

Again the complication introduced by  $u_1=1$  is avoided by decreasing  $t_1=0$ . The quantity  $g$  is a frequently occurring sum defined via

$$\sum_{i=1}^N \frac{\nabla s_i^2 \nabla s_i^2}{\nabla s_i} = 1 + g, \quad (3.13)$$

with explicit form

$$g = \sum_{i=1}^N s_i s_{i-1} (s_i - s_{i-1}) = \frac{1}{3}(1 - \delta g), \quad (3.14)$$

where

$$\delta g = \sum_{i=1}^N t_i^3. \quad (3.15)$$

Much of the mechanics of dealing with these sums have been worked out in Ref. [18]. However, their use and interpretation here are very different. From (2.4) and (2.5), we have

$$e_{TTV} = \frac{1}{12} + \frac{1}{2} g C_2, \quad (3.16)$$

$$e_{VTV} = \frac{1}{24} - \frac{1}{2} g C_2^2. \quad (3.17)$$

Both are now only functions of  $\{t_i\}$  through  $g$ .

### IV. FUNDAMENTAL THEOREMS

We can now prove a number of important results:

*Theorem 1.* For the class of extended-linear symplectic integrators defined by  $t_1=0$  and (3.8)–(3.10), if  $\{t_i\} > 0$  for  $i \neq 1$  such that  $e_T=1$ , then  $e_{TTV} \neq e_{VTV}$ .

*Proof.* Setting  $e_{TTV}=e_{VTV}$  produces a quadratic equation for  $C_2$ ,

$$C_2^2 + C_2 + \frac{1}{12g} = 0 \quad (4.1)$$

whose discriminant

$$D = b^2 - 4ac = -\frac{\delta g}{1 - \delta g} \quad (4.2)$$

is strictly negative (since if  $e_T=1$ , then  $1 > \delta g > 0$ ). Hence no real solution exists for  $C_2$ . This is a fundamental theorem about positive-coefficient factorizations. This was proved generally in the context of symplectic corrector (or process) algorithms by Chin [11] and by Blanes and Casas [19]. If  $e_{TTV}$  can never equal  $e_{VTV}$ , then no second-order algorithm with positive coefficients can be corrected beyond second order with the use of a corrector.

As a corollary, for  $\{t_{i>1}\} > 0$ ,  $e_{TTV}$  and  $e_{VTV}$  cannot both vanish. This is the content of the Sheng-Suzuki Theorem [20,21]: there are no integrators of an order greater than 2 of the form (2.1) with only positive factorization coefficients. Our proof here is restricted to extended-linear integrators, but can be interpreted more generally as it is done in Ref. [18]. Blanes and Casas [19] have also given elementary proof of this using a very weak *necessary* condition. Here, for extended-linear integrators, we can be very precise in stating how both  $e_{TTV}$  and  $e_{VTV}$  fail to vanish. We have, from (3.16), if  $e_{TTV}=0$ , then

$$C_2 = -\frac{1}{2(1 - \delta g)}, \quad e_{VTV} = -\frac{1}{24} \frac{\delta g}{(1 - \delta g)}. \quad (4.3)$$

Similarly, from (3.17), if  $e_{VTV}=0$ , then

$$C_2 = -\frac{1}{2\sqrt{1 - \delta g}}, \quad e_{TTV} = \frac{1}{12}(1 - \sqrt{1 - \delta g}). \quad (4.4)$$

Satisfying either condition forces  $C_2$  to be a function of  $\{t_i\}$  through  $\delta g$ . From Ref. [18], we have learned that the value given by (4.3) is actually an upperbound for  $e_{VTV}$  if  $\{t_{i>1}\} > 0$  and  $e_{TTV}=0$ . Similarly, in general, the value given by (4.4) is a lower bound for  $e_{TTV}$  if  $\{t_{i>1}\} > 0$  and  $e_{VTV}=0$ . Our class of extended-linear integrators are all algorithms that attain these bounds for positive  $t_{i>1}$ . Note that in (4.4) we have discarded the positive solution for  $C_2$  which would have led to negative values for the  $v_i$  coefficients.

For the study of forward integrators where one requires  $\{t_{i>1}\} > 0$ , it is useful to state (4.3) as a theorem:

*Theorem 2a.* For the class of extended-linear symplectic integrators defined by

$$v_1 = \frac{1}{2} + C_2(1 - t_2), \quad v_N = \frac{1}{2} + C_2(1 - t_N), \\ v_i = -C_2(t_i + t_{i+1}), \quad (4.5)$$

with  $t_1=0$ ,  $e_T=1$ , and  $C_2$ ,  $e_{VTV}$  given by

$$C_2 = -\frac{1}{2\phi}, \quad e_{VTV} = -\frac{1}{24} \left( \frac{1}{\phi} - 1 \right), \quad (4.6)$$

where

$$\phi = 1 - \delta g \quad \text{and} \quad \delta g = \sum_{i=1}^N t_i^3, \quad (4.7)$$

one has

$$\prod_{i=1}^N e^{t_i \varepsilon^T} e^{v_i \varepsilon^V} = \exp\{\varepsilon(T + V) + \varepsilon^3 e_{VTV}[V, [T, V]] + \dots\}. \quad (4.8)$$

For  $t_1=0$ , the first operator  $e^{v_1 \varepsilon^V}$  classically updates the velocity (momentum) variable. *Theorem 2a* completely described the structure of these *velocity*-type algorithms.

If one now interchanges  $T \leftrightarrow V$  and  $\{t_i\} \leftrightarrow \{v_i\}$ , then  $[T, [T, V]]$  transforms into  $[V, [T, V]]$  with a sign change. Hence, one needs to interpret  $e_{TTV}$  in (4.4) as  $-e_{VTV}$ , yielding as follows:

*Theorem 2b.* For the class of extended-linear symplectic integrators defined by

$$t_1 = \frac{1}{2} + C_2(1 - v_2), \quad t_N = \frac{1}{2} + C_2(1 - v_N), \\ t_i = -C_2(v_i + v_{i+1}), \quad (4.9)$$

with  $v_1=0$ ,  $e_V=1$ , and  $C_2$ ,  $e_{VTV}$  given by

$$C_2 = -\frac{1}{2\phi'}, \quad e_{VTV} = -\frac{1}{12}(1 - \phi'), \quad (4.10)$$

where

$$\phi' = \sqrt{1 - \delta g'} \quad \text{and} \quad \delta g' = \sum_{i=1}^N v_i^3, \quad (4.11)$$

one has

$$\prod_{i=1}^N e^{v_i \varepsilon^V} e^{t_i \varepsilon^T} = \exp\{\varepsilon(T + V) + \varepsilon^3 e_{VTV}[V, [T, V]] + \dots\}. \quad (4.12)$$

For  $v_1=0$ , the first operator  $e^{t_1 \varepsilon^T}$  classically updates the position variable. *Theorem 2b* completely described the structure of these *position*-type algorithms.

In both *Theorems 2a* and *2b*, one obtains fourth-order forward algorithms by simply moving the commutator  $[V, [T, V]]$  term back to the left-hand side and distribute it symmetrically among all the  $V$  operators [28].

If some  $t_i$  were allowed to be negative, then both  $e_{TTV}$  and  $e_{VTV}$  can be zero for  $\delta g=0$ . For both (4.3) and (4.4) we have

$$C_2 = -\frac{1}{2} \quad (4.13)$$

and

$$v_i = \frac{1}{2}(t_i + t_{i+1}). \quad (4.14)$$

The latter is now true even for  $i=1$  and  $i=N$ . This is not a coincident, from (3.16) and (3.17), if we set  $C_2=-\frac{1}{2}$ , then

$$e_{TTV} = 2e_{VTV} = \frac{1}{12} - \frac{g}{4} = \frac{1}{12} \delta g. \quad (4.15)$$

Since  $C_2$  here is a true constant,  $\{v_i\}$  is linearly related to  $\{t_i\}$ . We can formulate this explicitly as a theorem for the negative-coefficient factorization yielding truly linear algorithms:

*Theorem 3.* For the class of truly *linear* algorithms defined by

$$v_1 = \frac{1}{2}t_2, \quad v_N = \frac{1}{2}t_N, \quad v_i = \frac{1}{2}(t_i + t_{i+1}), \quad (4.16)$$

and  $t_1=0$ , one has

$$\prod_{i=1}^N e^{t_i \varepsilon T} e^{v_i \varepsilon V} = \exp \left\{ \varepsilon e_T (T + V) + \frac{\varepsilon^3}{24} \delta g (2[T, [T, V]] + [V, [T, V]]) + \dots \right\}. \quad (4.17)$$

Both commutators now vanish simultaneously if  $\delta g=0$ .

*Theorem 3* can be proven more directly by noting that

$$\mathcal{T}_2(t_i) = e^{1/2 t_i \varepsilon V} e^{t_i \varepsilon T} e^{1/2 t_i \varepsilon V} = \exp \left[ t_i \varepsilon (T + V) + t_i^3 \frac{\varepsilon^3}{24} (2[T, [T, V]] + [V, [T, V]]) + O(\varepsilon^5) \right], \quad (4.18)$$

the product  $\prod_{i=2}^N \mathcal{T}_2(t_i)$  then reproduces (4.17). This has been derived by Blanes and Casas [19] in a discussion before their *Theorem 5*. However, they were more interested in using *Theorem 3* above to discuss the distribution of negative coefficients than in deriving fourth-order algorithms. For example, an immediate corollary of *Theorem 3* is that if  $\delta g$  were to vanish, then there must be at least one  $t_k < 0$  such that  $t_k^3 + t_{k+1}^3 < 0$  or  $t_k^3 + t_{k-1}^3 < 0$ . Since

$$(x^3 + y^3) = (x + y) \left[ \frac{3}{4} y^2 + \left( x - \frac{1}{2} y \right)^2 \right],$$

$x^3 + y^3 < 0 \Rightarrow x + y < 0$ . We therefore must have  $t_k + t_{k+1} < 0$  or  $t_k + t_{k-1} < 0$ . From (4.16), this implies that  $v_k$  or  $v_{k-1}$  must be negative. Thus an algorithm of order greater than 2 of the form (4.17) must contain at least one pair of negative coefficients  $t_i$  and  $v_j$ . In its general context, this is the Goldman-Kaper result [22]. Our linear formulation here, as well as Blanes and Casas' *Theorem 5*, is more precise: if  $t_i$  is negative, then at least one of its adjacent  $v_i$  must be negative. If only one  $t_k$  is negative, then both of its adjacent  $v_i$  must be negative. For further discussions on the distribution of negative coefficients, see Blanes and Casas [19].

## V. THE STRUCTURE OF FORWARD INTEGRATORS

*Theorems 2a* and *2b* can be used to construct fourth-order forward algorithms with only positive factorization coefficients. These forward integrators are the only fourth-order factorized symplectic algorithms capable of integrating *time-irreversible* equations such as the Fokker-Planck [10,23] or the imaginary time Schrödinger equation [24–26]. Since it has been shown that [18] currently there are no practical ways of constructing sixth-order forward integrators, these fourth-order algorithms enjoy a unique status.

For  $N=3$ , for a fourth-order algorithm, we must require  $t_2=t_3=\frac{1}{2}$ . *Theorem 2a* then implies that

$$v_1 = v_3 = \frac{1}{6}, \quad v_2 = \frac{2}{3}, \quad \text{and } e_{VTV} = -\frac{1}{72}. \quad (5.1)$$

By moving the term  $\varepsilon^3 e_{VTV} [V, [T, V]]$  back to the left-hand side of (1.2) and combining it with the central  $V$ , one recov-

ers forward algorithm 4A [27,28]. For  $N=4$  with  $t_2=t_3=t_4=\frac{1}{3}$ , we have

$$v_1 = v_4 = \frac{1}{8}, \quad v_2 = v_3 = \frac{3}{8}, \quad \text{and } e_{VTV} = -\frac{1}{192}, \quad (5.2)$$

which corresponds to forward algorithm 4D [13]. These are special cases of the general *minimal*  $|e_{VTV}|$ , velocity-type algorithm given by  $t_1=0$ ,  $t_i=1/(N-1)$ ,

$$v_1 = v_N = \frac{1}{2N}, \quad v_i = \frac{(N-1)}{N(N-2)}, \quad \text{with } e_{VTV} = -\frac{1}{24} \frac{1}{N(N-2)}. \quad (5.3)$$

This arbitrary  $N$  algorithm can serve as a useful check for any general fourth-order, velocity-type algorithm.

Alternatively, for  $N=4$ , we can allow  $t_2$  to be a free parameter so that

$$t_4 = t_2, \quad t_3 = 1 - 2t_2. \quad (5.4)$$

*Theorem 2a* then fixes  $C_2$  and  $e_{VTV}$  with

$$\phi = 6t_2(1 - t_2)^2 \quad (5.5)$$

and

$$v_2 = v_3 = \frac{1}{12t_2(1 - t_2)}, \quad v_1 = v_4 = \frac{1}{2} - v_2. \quad (5.6)$$

One recognizes that this is the one-parameter algorithm 4BDA first found in Ref. [14] using symbolic algebra. For  $t_2=\frac{1}{2}$ , one recovers the integrator 4A; for  $t_2=\frac{1}{3}$ , one gets back 4D. The advantage of using a variable  $t_2$  is that one can use it to minimize the resulting fourth-order error (oftentime to zero) in any specific application. All the seven-stage, forward integrators in the velocity form described by Omelyan, Mryglod, and Folk (OMF) [17] correspond to different ways of choosing  $t_2$  and distributing the commutator term in 4BDA.

For  $N=5$ , again using  $t_2$  as a parameter, we have  $t_1=0$ ,  $t_5=t_2$ ,  $t_4=t_3=\frac{1}{2}-t_2$ , (4.6) with

$$\phi = \frac{15}{16} - 3\left(t_2 - \frac{1}{4}\right)^2, \quad (5.7)$$

$v_5=v_1$ ,  $v_4=v_2$ ,  $v_3=1-2(v_1+v_2)$ , and

$$v_1 = \frac{1}{2} + C_2(1 - t_2), \quad v_2 = -\frac{1}{2}C_2. \quad (5.8)$$

This is a new one-parameter family of fourth-order algorithms with nine stages or operators.

To generate position-type algorithms, one can apply *Theorem 2b*. For  $N=3$ , with  $v_1=0$ ,  $v_1=v_2=\frac{1}{2}$ , we have

$$t_1 = t_3 = \frac{1}{2} \left( 1 - \frac{1}{\sqrt{3}} \right), \quad t_2 = \frac{1}{\sqrt{3}}, \quad \text{and } e_{VTV} = -\frac{1}{12} \left( 1 - \frac{1}{2}\sqrt{3} \right). \quad (5.9)$$

This produces forward algorithm 4B [27,28] corresponding to  $t_2=(1-1/\sqrt{3})/2$  in 4BDA. Again, this is a special case of the general fourth-order, minimal  $|e_{VTV}|$  algorithm with  $v_1=0$ ,  $v_i=1/(N-1)$ ,

$$t_1 = t_N = \frac{1}{2} \left( 1 - \sqrt{\frac{N-2}{N}} \right), \quad t_i = \frac{1}{\sqrt{N(N-2)}}, \quad (5.10)$$

and

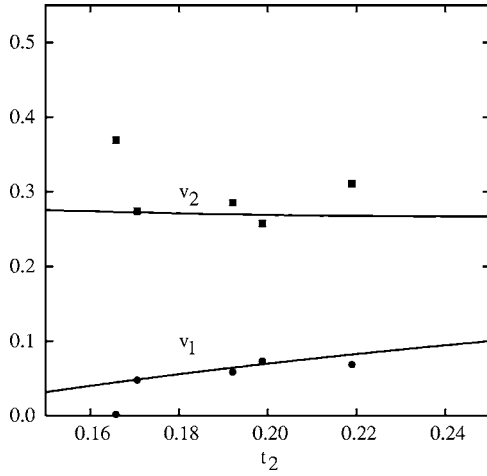


FIG. 1. Comparing the coefficients of five, 9-stage, *velocity*-type, fourth-order forward integrators of Omelyan, Mryglod, and Folk [17] (filled circles and squares), with the analytical prediction of extended-linear symplectic integrators (solid lines).

$$e_{vTV} = -\frac{1}{12} \left( 1 - \frac{\sqrt{N(N-2)}}{(N-1)} \right). \quad (5.11)$$

For  $N=4$ ,  $v_1=0$  and  $v_2$  as the free parameter, invoking *Theorem 2b* gives

$$v_4 = v_2, \quad v_3 = 1 - 2v_2, \quad (5.12)$$

$$t_2 = t_3 = \frac{1}{2\sqrt{6v_2}}, \quad t_1 = t_4 = \frac{1}{2} - t_2 \quad (5.13)$$

and

$$e_{vTV} = -\frac{1}{12} [1 - (1 - v_2)\sqrt{6v_2}]. \quad (5.14)$$

For  $v_2 = \frac{1}{6}$  and  $v_2 = \frac{3}{8}$ , this reproduces algorithms 4A and 4C [28], respectively. One again recognizes that the above is the one-parameter algorithm 4ACB first derived in Ref. [14], but now with a much simpler parametrization. Algorithm 4ACB covers all of the seven-stage, forward fourth-order position-type integrators described by OMF [17].

For  $N=5$ , with  $v_2$  as a free parameter, we have  $v_1=0$ ,  $v_5=v_2$ ,  $v_4=v_3 = \frac{1}{2} - v_2$ , and *Theorem 2b* produces another 9-stage fourth-order algorithm with

$$\phi' = \sqrt{\frac{15}{16} - 3\left(v_2 - \frac{1}{4}\right)^2}. \quad (5.15)$$

$t_5=t_1$ ,  $t_4=t_2$ ,  $t_3=1-2(t_1+t_2)$ , and

$$t_1 = \frac{1}{2} + C_2(1 - v_2), \quad t_2 = -\frac{1}{2}C_2. \quad (5.16)$$

For  $N < 5$ , we have shown above that all fourth-order algorithms are necessarily extended-linear. For  $N \geq 5$ , this is not necessarily the case. Nevertheless we find that, remarkably, most known  $N=5$  (9 stages) forward algorithms are very close to being extended-linear. For velocity-type,  $N=5$  extended-linear algorithms,  $v_1$  and  $v_2$  are functions of  $t_2$  fixed by (5.8). In Fig. 1, we compare this predicted relationship with the actual values of  $v_1$ ,  $v_2$ , and  $t_2$  of five forward, velocity-type, fourth-order algorithms found by OMF [17].

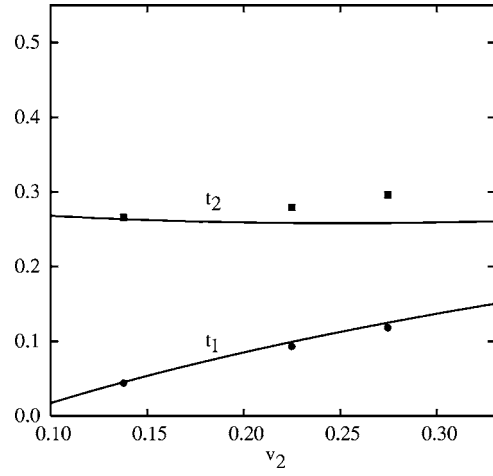


FIG. 2. Comparing the coefficients of three, 9-stage, *position*-type, fourth-order forward integrators of Omelyan, Mryglod, and Folk [17] (filled circles and squares), with the analytical prediction of extended-linear symplectic integrators (solid lines).

These are their Eqs. (52)–(56), and with their  $\theta$ ,  $\vartheta$ , and  $\lambda$  correspond to  $t_2$ ,  $v_1$ , and  $v_2$ , respectively. Four of their five algorithms, with  $v_1$  in particular, are well described by (5.8).

In Fig. 2, we compare the coefficients of all three of OMF's forward, position-type algorithms, Eqs. (59)–(61), with (5.16) which fixes  $t_1$ ,  $t_2$  as a function of  $v_2$ . Here, their parameters  $\lambda$ ,  $\rho$ ,  $\theta$  correspond to  $v_2$ ,  $t_1$ ,  $t_2$ , respectively. Again,  $t_1$  is particularly well predicted by (5.16).

For 11-stage algorithms with  $N=6$ , we have two free parameters  $t_2$ ,  $t_3$  for velocity-type algorithms with

$$\phi = 1 - 2t_2^3 - 2t_3^3 - (1 - 2t_2 - 2t_3)^3 \quad (5.17)$$

and two free parameters  $v_2, v_3$  for position-type algorithms with

$$\phi' = \sqrt{1 - 2v_2^3 - 2v_3^3 - (1 - 2v_2 - 2v_3)^3}. \quad (5.18)$$

Once  $\phi$  and  $\phi'$  are known, we can determine  $v_1$  and  $v_2$  in the case of velocity-type algorithms and  $t_1$  and  $t_2$  in the case of position-type algorithms. There is one 11-stage velocity algorithm with positive coefficients found by OMF; their Eq. (68) with  $\rho(=t_2)=0.2029$ ,  $\theta(=t_3)=0.1926$ ,

$$\vartheta(=v_1) = 0.0667, \quad \text{and} \quad \lambda(=v_2) = 0.2620. \quad (5.19)$$

The last two values are to be compared with the values given by *Theorem 2a* below at the same values of  $t_2$  and  $t_3$ ,

$$v_1 = 0.0848, \quad \text{and} \quad v_2 = 0.2060. \quad (5.20)$$

For OMF's 11-stage, position-type algorithm Eq. (78), with  $\vartheta(=v_2)=0.1518$ ,  $\lambda(=v_3)=0.2158$ ,

$$\rho(=t_1) = 0.0642, \quad \text{and} \quad \theta(=t_2) = 0.1920. \quad (5.21)$$

For the same values of  $v_2$  and  $v_3$ , *Theorem 2b* gives

$$t_1 = 0.0659, \quad \text{and} \quad t_2 = 0.1881. \quad (5.22)$$

It is remarkable that these 11-stage, fourth-order algorithms derived by complex symbolic algebra, remained very close to the values predicted by our extended-linear algorithms.

## VI. THE STRUCTURE OF NONFORWARD INTEGRATORS

*Theorem 3* can be used to derive two distinct families of nonforward, fourth-order algorithms. Consider first the case of  $N=4$ . For  $t_1=0$  with symmetric coefficients  $t_4=t_2$ , the constraints

$$2t_2 + t_3 = 1, \quad (6.1)$$

$$2t_2^3 + t_3^3 = 0, \quad (6.2)$$

have unique solutions

$$t_2 = \frac{1}{2 - 2^{1/3}} \text{ and } t_3 = -\frac{2^{1/3}}{2 - 2^{1/3}}. \quad (6.3)$$

Equation (4.16) then yields

$$v_1 = v_4 = \frac{1}{2} \frac{1}{2 - 2^{1/3}}, \quad v_2 = v_3 = -\frac{1}{2} \frac{(2^{1/3} - 1)}{2 - 2^{1/3}}. \quad (6.4)$$

One recognizes that we have just derived the well-known fourth-order Forest-Ruth integrator [29]. Note that there is complete symmetry between  $\{t_i\}$  and  $\{v_i\}$ . For position-type algorithms, we simply interchange the values of  $t_i$  and  $v_i$ .

There are no symmetric solutions for  $N=5$ , for the same reason that there are also no solutions for  $N=3$ . For  $N=2k$ , we have the general condition

$$\begin{aligned} 2 \sum_{i=2}^k t_i + t_{k+1} &= 1, \\ 2 \sum_{i=2}^k t_i^3 + t_{k+1}^3 &= 0, \end{aligned} \quad (6.5)$$

which can be solved by introducing real parameters  $\alpha_i$  for  $i=2$  to  $k$  with  $\alpha_2=1$ ,

$$t_i = \alpha_i t_2, \quad (6.6)$$

so that

$$t_{k+1} = -2^{1/3} \left( \sum_{i=2}^k \alpha_i^3 \right)^{1/3} t_2, \quad (6.7)$$

$$t_2 = \frac{1}{2 \left( \sum_{i=2}^k \alpha_i \right) - 2^{1/3} \left( \sum_{i=2}^k \alpha_i^3 \right)^{1/3}}. \quad (6.8)$$

These solutions generalize the fourth-order Forest-Ruth integrator to arbitrary  $N$ .

The general fourth-order condition (6.5) has been derived previously by McLachlan [30] using the generalized triplet

construction published by Creutz and Gocksch [31]. However, the invocation of *Theorem 3* is more general and much simpler. McLachlan suggests that one should just set all  $\alpha_i = 1$ .

For  $N=2k+1$ ,  $k>2$ , again introducing (6.6) for  $i=2$  to  $k$  with  $\alpha_2=1$ , we have

$$t_{k+1} = - \left( \sum_{i=2}^k \alpha_i^3 \right)^{1/3} t_2, \quad (6.9)$$

$$t_2 = \frac{1}{2 \left( \sum_{i=2}^k \alpha_i \right) - 2 \left( \sum_{i=2}^k \alpha_i^3 \right)^{1/3}}. \quad (6.10)$$

This is a new class of a fourth-order algorithm possible only for  $N$  odd greater than 5 and is not derivable from the triplet construction.

## VII. CONCLUSIONS

Most of the machinery for tracking coefficients was developed in Ref. [18] for the purpose of providing a constructive proof of the Sheng-Suzuki theorem. The advantage of this constructive approach is that we can obtain explicit lower bounds on the second-order error coefficients. Here, by imposing the extended-linear relationship between  $\{t_i\}$  and  $\{v_i\}$ , these bounds become the actual error coefficients and provide a complete characterization for all fourth-order symplectic integrators for an arbitrary number of operators. The most satisfying aspect of this work is that most fourth-order integrators can now be derived analytically without recourse to symbolic algebra or numerical root finding. We have also provided explicit construction of many new classes of fourth-order algorithms.

For  $N=5, 6$ , corresponding to 9 and 11 operators, we have shown that many fourth-order algorithms found by Omelyan, Mryglod, and Folk [17] are surprisingly close to the predicted coefficients of our theory, suggesting that the extended-linear relation between coefficients may be the dominate solution of the order condition.

The expansion (3.2) may hold similar promise for characterizing sixth-order algorithms by introducing extended-quadratic or higher order relationships between the two sets of coefficients.

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- [1] H. Yoshida, *Celest. Mech.* **56**, 27 (1993).
- [2] E. Hairer, C. Lubich, and G. Wanner, *Geometric Numerical Integration* (Springer-Verlag, Berlin, 2002).
- [3] R. I. McLachlan and G. R. W. Quispel, *Acta Numerica* **11**, 241 (2002).
- [4] S. A. Chin and C. R. Chen, *Celest. Mech. Dyn. Astron.* **91**, 301 (2005).
- [5] S. R. Scuro and S. A. Chin, *Phys. Rev. E* **71**, 056703 (2005).
- [6] T. Hirono, W. Lui, S. Seki, and Y. Yoshikuni, *IEEE Trans. Microwave Theory Tech.* **49**, 1640 (2001).
- [7] J. Lee and B. Fornberg, *J. Comput. Appl. Math.* **158**, 485 (2003).
- [8] M. Takahashi and M. Imada, *J. Phys. Soc. Jpn.* **53**, 3765 (1984).
- [9] S. Jang, S. Jang, and G. A. Voth, *J. Chem. Phys.* **115**, 7832 (2001).
- [10] H. A. Forbert and S. A. Chin, *Phys. Rev. E* **63**, 016703 (2001).
- [11] S. A. Chin, *Phys. Rev. E* **69**, 046118 (2004).
- [12] D. Feit, J. A. Fleck, Jr., and A. Steiger, *J. Comput. Phys.* **47**, 412 (1982).
- [13] S. A. Chin and C. R. Chen, *J. Chem. Phys.* **114**, 7338 (2001).
- [14] S. A. Chin and C. R. Chin, *J. Chem. Phys.* **117**, 1409 (2002).
- [15] P. V. Koseleff, *Integration Algorithms and Classical Mechanics*, Fields Institute Communications, Vol. 10 (American Mathematical Society, Providence, RI, 1996), p. 103.
- [16] I. P. Omelyan, I. M. Mryglod, and R. Folk, *Phys. Rev. E* **66**, 026701 (2002).
- [17] I. P. Omelyan, I. M. Mryglod, and R. Folk, *Comput. Phys. Commun.* **151**, 272 (2003).
- [18] S. A. Chin, *Phys. Rev. E* **71**, 016703 (2005).
- [19] S. Blanes and F. Casas, *Appl. Numer. Math.* **54**, 23 (2005).
- [20] Q. Sheng, *IMA J. Numer. Anal.* **9**, 199 (1989).
- [21] M. Suzuki, *J. Math. Phys.* **32**, 400 (1991).
- [22] D. Goldman and T. J. Kaper, *IMA J. Numer. Anal.* **33**, 349 (1996).
- [23] H. A. Forbert and S. A. Chin, *Phys. Rev. B* **63**, 144518 (2001).
- [24] J. Auer, E. Krotscheck, and S. A. Chin, *J. Chem. Phys.* **115**, 6841 (2001).
- [25] O. Ciftja and S. A. Chin, *Phys. Rev. B* **68**, 134510 (2003).
- [26] S. A. Chin and E. Krotscheck, *Phys. Rev. E* **72**, 036705 (2005).
- [27] M. Suzuki, in *Computer Simulation Studies in Condensed Matter Physics VIII*, edited by D. Landau, K. Mon, and H. Shuttler (Springler, Berlin, 1996).
- [28] S. A. Chin, *Phys. Lett. A* **226**, 344 (1997).
- [29] E. Forest and R. D. Ruth, *Physica D* **43**, 105 (1990).
- [30] R. I. McLachlan, *Numer. Algorithms* **31**, 233 (2002).
- [31] M. Creutz and A. Gocksch, *Phys. Rev. Lett.* **63**, 9 (1989).