

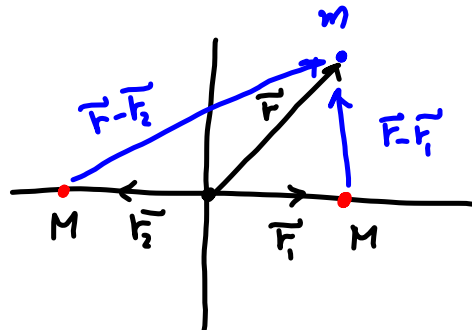
Solving for explicitly time-dependent force problems: $\vec{a}(\vec{r}, t)$

Gravitational 3-body problem \leftrightarrow no analytical solution

Sun-earth-moon

\hookrightarrow most orbits are chaotic!

2) Two-center problem - two gravitational attraction centers.



$$m\ddot{\vec{r}} = -GMm \frac{\vec{r}-\vec{r}_1}{|\vec{r}-\vec{r}_1|^3} - GMm \frac{\vec{r}-\vec{r}_2}{|\vec{r}-\vec{r}_2|^3}$$

choose:

$$|\vec{r}_1| = |\vec{r}_2| = \frac{1}{2}$$

$$\vec{r}_1 = (\frac{1}{2}, 0)$$

$$\vec{r}_2 = (-\frac{1}{2}, 0)$$

$$\ddot{\vec{r}} = -\frac{1}{2} \left(\frac{\vec{d}_1}{|\vec{d}_1|^3} + \frac{\vec{d}_2}{|\vec{d}_2|^3} \right)$$

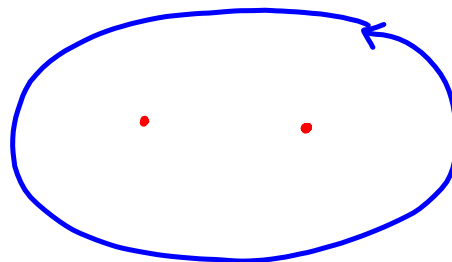
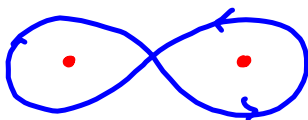
where $\vec{d}_1 = \vec{r} - \vec{r}_1$, $\vec{d}_2 = \vec{r} - \vec{r}_2$

Possible orbits



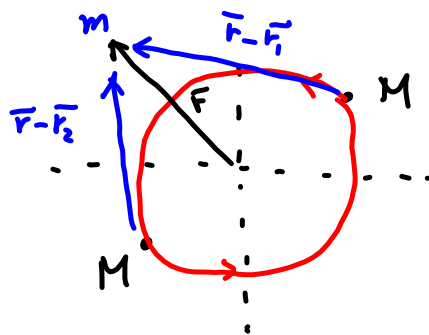
3)

2)



new figure 8-orbit

The **circular-restricted** 3-body problem



$$m \ll M$$

Assume that m has no effect on the motions of M_s .

$$\vec{r}_1(t) = \frac{1}{2} (\cos(\omega t), \sin(\omega t))$$

$$\vec{r}_2(t) = -\vec{r}_1(t)$$

$$\omega = 1$$

$$\ddot{\vec{r}} = \vec{a}(\vec{r}, t)$$

$$= -\frac{1}{2} \left(\frac{\vec{d}_1(t)}{|\vec{d}_1(t)|^3} + \frac{\vec{d}_2(t)}{|\vec{d}_2(t)|^3} \right)$$

where $\vec{d}_1(t) = \vec{r} - \vec{r}_1(t)$, $\vec{d}_2 = \vec{r} - \vec{r}_2(t)$

\Rightarrow motion of m is a time-varying gravitational field.

Symplectic algorithms for solving
a time-dependent $\vec{a}(\vec{r}, t)$

For the case sym 1a + 1b algorithm

1a) Cromer

1) $\vec{v}_1 = \vec{v}_0 + \vec{a}(\vec{r}_0, t) \Delta t$ $t=0$

2) $\vec{r}_1 = \vec{r}_0 + \vec{v}_1 \Delta t$

3) $t = t + \Delta t$

then repeat

for-loop initial time

evaluate \vec{a} at the initial time

1b)

1) $\vec{r}_1 = \vec{r}_0 + \vec{v}_0 \Delta t$

2) $t = t + \Delta t$

3) $\vec{v}_1 = \vec{v}_0 + \vec{a}(\vec{r}_1, t)$

for-loop

evaluate at the end-time.

Second-order 2b (don't use 2a)
 ↑ position - $\Delta x / \Delta t$ ↑ velocity - $\Delta v / \Delta t$

2b)

$$\vec{r}_{2b} = \vec{r}_{t=0} + \frac{1}{2} \Delta t \vec{v} + \vec{v} \Delta t + \frac{1}{2} \Delta t \vec{a}$$

Suzuki's Rule:

update t
 after every
 updating of
 \vec{r} , with
 the coefficient of
 v .

$$\vec{r}_1 = \vec{r}_0 + \frac{1}{2} \Delta t \vec{v}_0$$

$t = t + \frac{1}{2} \Delta t$

$$\vec{v}_1 = \vec{v}_0 + \Delta t \vec{a}(\vec{r}_1, t)$$

$$\vec{r}_2 = \vec{r}_1 + \frac{1}{2} \Delta t \vec{v}_1$$

$t = t + \frac{1}{2} \Delta t$

} for loop

Algorithm 2a: Velocity-Verlet

$$T_{2a} = e^{\frac{1}{2}\Delta t V} e^{\Delta t T} e^{\frac{1}{2}\Delta t V}$$

$t=0$

Algorithm

$$\bar{v}_1 = \bar{v}_0 + \frac{1}{2}\Delta t \bar{a}(\bar{r}_0, t)$$

$$\bar{r}_1 = \bar{r}_0 + \Delta t \bar{v}_1$$

$$\bar{v}_2 = \bar{v}_1 + \frac{1}{2}\Delta t \bar{a}(\bar{r}_1, t)$$

update t
after each position
update, with
the coefficient
of v .

for-loop

Solve time-dependent Hamiltonians

$$H(t) = T + V(t)$$

Recall previously $\frac{dW}{dt} = HW \rightarrow \frac{dW}{W} = H dt$

$$\Rightarrow \boxed{W(t) = e^{\int_0^t dt' H(t')} W(0)} \quad \begin{array}{l} d \ln W = H dt \\ \ln W = \int H(t) dt + C \end{array}$$

true only for $H(t; \mathbf{r}, \mathbf{p})$ a fct, no operators.

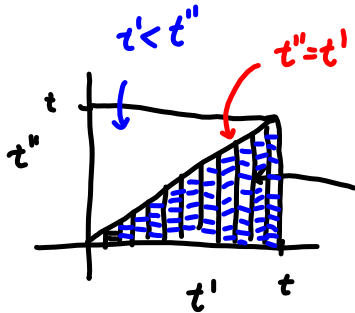
Solve $\frac{dw}{dt} = \hat{H}(t)w$ $\hat{H} = \text{operator}$
 $= T + V(x)$

$w(t) = u(t)w(0)$

evolution operator $\int_0^t \frac{du}{dt} w(0) = \int_0^t \hat{H}(t') u(t') w(0) dt'$ $\sum_i v_i \frac{\partial}{\partial q_i}$
 $\sum_i a(q_i, t) \frac{\partial}{\partial v}$

$u(t) - u(0) = \int_0^t \hat{H}(t') u(t') dt'$
 " 1

$u(t) = 1 + \int_0^t \hat{H}(t') u(t') dt'$



$= 1 + \int_0^t \hat{H}(t') [1 + \int_0^{t'} \hat{H}(t'') u(t'') dt''] dt'$
 $= 1 + \int_0^t \hat{H}(t') dt' + \int_0^t dt' \int_0^{t'} dt'' \hat{H}(t') \hat{H}(t'')$
 $\int_0^t dt'' \int_0^{t''} dt' \hat{H}(t'') \hat{H}(t')$
 " " $t' < t''$ $t' > t''$

Define: time-ordering of operators: $\hat{H}(t') \hat{H}(t'')$
 " " $= \hat{H}(t'') \hat{H}(t')$

$T[\hat{H}(t') \hat{H}(t'')] [\hat{H}(t'), \hat{H}(t'')] = 0$
 $= \begin{cases} \hat{H}(t') \hat{H}(t'') & \text{when } t' > t'' \text{ if true} \\ \hat{H}(t'') \hat{H}(t') & \text{when } t' < t'' \end{cases}$
 but $[T+V(t'), T+V(t'')] \neq 0$

$= 1 + \frac{1}{2} \int_0^t dt'' \int_0^{t''} dt' \hat{H}(t'') \hat{H}(t')$
 $= 1 + \frac{1}{2} \left[\int_0^t dt' \hat{H}(t') \right]^2 + \dots$
 $= \int_0^t dt' \hat{H}(t')$
 $= 1 + \frac{1}{2} \int_0^t dt'' \int_0^{t''} dt'$
 $T[\hat{H}(t') \hat{H}(t'')]$

$\left[v \frac{\partial}{\partial q} + a(t') \frac{\partial}{\partial v} \right] \left[v \frac{\partial}{\partial q} + a(t'') \frac{\partial}{\partial v} \right]$
 $v \frac{\partial}{\partial q} \quad v \frac{\partial}{\partial q} \quad a(t') \frac{\partial}{\partial v} \quad a(t'') \frac{\partial}{\partial v}$

with time-ordering

$$U(t) = 1 + \int_0^t \hat{H}(t') dt' + \frac{1}{2} \int_0^t dt' \int_0^{t'} dt'' T[\hat{H}(t') \hat{H}(t'')] + \frac{1}{3!} \int_0^t dt_1 \int_0^{t_1} dt_2 \int_0^{t_2} dt_3 T[\hat{H}(t_1) \hat{H}(t_2) \hat{H}(t_3)] + \dots$$

$$U(t) = T \left(e^{\int_0^t dt' \hat{H}(t')} \right) \leftarrow \text{seems difficult to implement}$$

M. Suzuki
JCP 117 (2002) 1403
↑ Japanese

$$= T \left(e^{\Delta t \sum_{k=0}^n \hat{H}(k\Delta t)} \right)$$

$$= e^{\Delta t \hat{H}(t)} e^{\Delta t \hat{H}(t-\Delta t)} \dots e^{\Delta t \hat{H}(0)}$$

↑ manifestly time-ordered!